

# Rotational Motion

Reading: Chapter 10

## Angular Displacement

$$\Delta\theta = \theta_2 - \theta_1.$$

## Angular Velocity

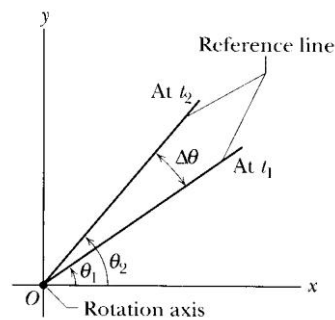


FIG. 10-4 The reference line of the rigid body of Figs. 10-2 and 10-3 is at angular position  $\theta_1$  at time  $t_1$  and at angular position  $\theta_2$  at a later time  $t_2$ . The quantity  $\Delta\theta (= \theta_2 - \theta_1)$  is the angular displacement that occurs during the interval  $\Delta t (= t_2 - t_1)$ . The body itself is not shown.

## Average angular velocity

$$\bar{\omega} = \frac{\theta_2 - \theta_1}{t_2 - t_1} = \frac{\Delta\theta}{\Delta t}.$$

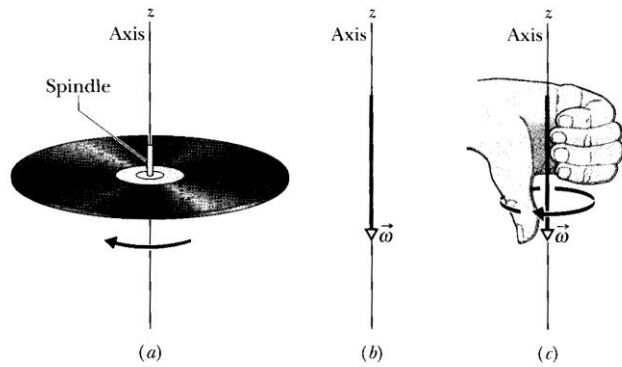
## Instantaneous angular velocity

$$\omega = \lim_{\Delta t \rightarrow 0} \frac{\Delta\theta}{\Delta t} = \frac{d\theta}{dt}.$$

The magnitude is called the **angular speed**.

## Angular Velocity as a Vector

**FIG. 10-6** (a) A record rotating about a vertical axis that coincides with the axis of the spindle. (b) The angular velocity of the rotating record can be represented by the vector  $\vec{\omega}$ , lying along the axis and pointing down, as shown. (c) We establish the direction of the angular velocity vector as downward by using a right-hand rule. When the fingers of the right hand curl around the record and point the way it is moving, the extended thumb points in the direction of  $\vec{\omega}$ .



The direction of the vector  $\vec{\omega}$  points along the axis of rotation, according to the **right-hand rule**.

## Angular Acceleration

Average angular acceleration

$$\bar{\alpha} = \frac{\omega_2 - \omega_1}{t_2 - t_1} = \frac{\Delta\omega}{\Delta t}.$$

Instantaneous angular acceleration

$$\alpha = \lim_{\Delta t \rightarrow 0} \frac{\Delta\omega}{\Delta t} = \frac{d\omega}{dt}.$$

## Rotation with Constant Acceleration

$$\omega = \omega_0 + \alpha t,$$

$$\theta = \omega_0 t + \frac{1}{2} \alpha t^2,$$

$$\omega^2 = \omega_0^2 + 2\alpha\theta.$$

## Example

**10-3** A grindstone rotates at constant angular acceleration  $\alpha = 0.35 \text{ rad s}^{-2}$ . At time  $t = 0$ , it has angular velocity  $\omega_0 = -4.6 \text{ rad s}^{-1}$  and a reference line on it is horizontal, at the angular position  $\theta_0 = 0$ .

(a) At what time after  $t = 0$  is the reference line at the angular position  $\theta = 5 \text{ rev}$ ?

(b) Describe the grindstone's rotation between  $t = 0$  and  $t = 32 \text{ s}$ .

(c) At what time  $t$  does the grindstone momentarily stop?

(a) Since  $\theta = \omega_0 t + \frac{1}{2} \alpha t^2$ ,

$$10\pi = -4.6t + \left(\frac{1}{2}\right)(0.35)t^2$$

$$0.35t^2 - 9.2t - 20\pi = 0$$

$$t = \frac{9.2 \pm \sqrt{9.2^2 + (4)(0.35)(20\pi)}}{0.7} = 31.9 \text{ or } -5.63$$

Therefore,  $t = 32 \text{ s}$ . (ans)

(b) The wheel is initially rotating in the clockwise direction. It slows down, stops, and then reverses to rotate in the anticlockwise direction. The reference line returns to its initial orientation of  $\theta = 0$ , and turns an additional 5 rev by time  $t = 32 \text{ s}$ . (ans)

(c) Since  $\omega = \omega_0 + \alpha t$ ,

$$0 = -4.6 + 0.35t$$

$$t = \frac{0 - (-4.6)}{0.35} = 13 \text{ s} \quad (\text{ans})$$

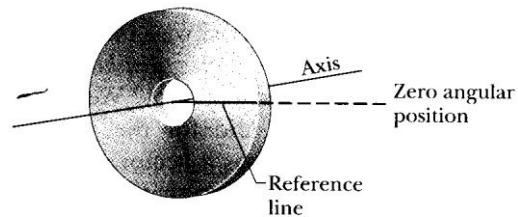


FIG. 10-8 A grindstone. At  $t = 0$  the reference line (which we imagine to be marked on the stone) is horizontal.

## Relating the Linear and Angular Variables

The position

$$s = \theta r.$$

The velocity

$$v = \omega r.$$

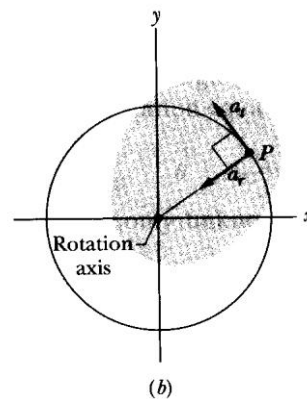
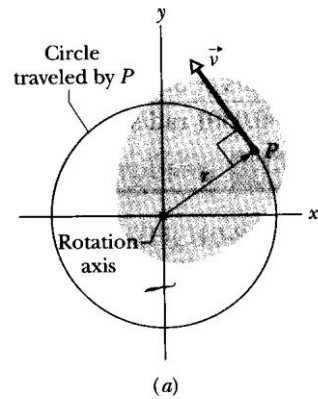
The acceleration

Tangential component:

$$a_t = \alpha r.$$

Radial component:

$$a_r = \frac{v^2}{r} = \omega^2 r.$$



**FIG. 10-9** The rotating rigid body of Fig. 10-2, shown in cross section viewed from above. Every point of the body (such as  $P$ ) moves in a circle around the rotation axis. (a) The linear velocity  $\vec{v}$  of every point is tangent to the circle in which the point moves. (b) The linear acceleration  $\vec{a}$  of the point has (in general) two components: tangential  $a_t$  and radial  $a_r$ .

## Example

**10-5** A roller coaster track is designed as follows.

(1) The passenger leaves the loading point with acceleration  $g$  along the horizontal track.

(2) The first section of the track forms a circular arc, so that the passenger also experiences a centripetal acceleration.

(3) When the magnitude  $a$  of the net acceleration reaches  $4g$  at some point  $P$  and angle  $\theta_P$  along the arc, the passenger then moves in a straight line along the tangent of the arc.

(a) What is the angle  $\theta_P$ ?

(b) What is the magnitude  $a$  of the passenger's net acceleration at  $P$  and after  $P$ ?

(a) Tangential acceleration:

$$a_t = g$$

Centripetal acceleration:  $a_r$   
 $= r\omega^2$

Since

$$\omega^2 = \omega_0^2 + 2\alpha\theta = 2\alpha\theta \quad \text{and}$$

$$\alpha = \frac{a_t}{r} = \frac{g}{r}, \quad \text{we have} \quad a_r = r(2\alpha\theta) = r\left(\frac{2g\theta}{r}\right) = 2g\theta$$

Magnitude of acceleration:  $a = \sqrt{a_t^2 + a_r^2} = \sqrt{g^2 + (2g\theta)^2}$   
 $= g\sqrt{1 + 4\theta^2}$

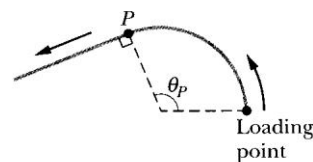
When  $a$  reaches  $4g$  at  $P$ ,  $4g = g\sqrt{1 + 4\theta^2}$

$$\Rightarrow \theta = \sqrt{\frac{15}{4}} = 1.94 \text{ rad} = 111^\circ \quad (\text{ans})$$

(b) At  $P$ ,  $a = 4g$ . (ans)

After  $P$ ,  $a_t = g$  and  $a_r = 0$ . Therefore,  $a = g$ . (ans)

Remark: This abrupt change in acceleration can cause roller-coaster headache.



**10-10** An overhead view of a horizontal track for a roller coaster. The track begins as a circular arc at the loading point and then, at point  $P$ , continues along a tangent to the arc.

## Kinetic Energy of Rotation

Consider the kinetic energy of a rotating rigid body:

$$K = \frac{1}{2}m_1v_1^2 + \frac{1}{2}m_2v_2^2 + \dots = \sum_i \frac{1}{2}m_iv_i^2.$$

Since  $v = \omega r$ , and  $\omega$  is the same for all particles, we have

$$K = \sum_i \frac{1}{2}m_i(\omega r_i)^2 = \frac{1}{2} \left( \sum_i m_i r_i^2 \right) \omega^2.$$

$\sum_i m_i r_i^2$  is called the **rotational inertia**. It tells us how the mass of the rotating body is distributed about its axis of rotation. In summary,

$$\boxed{I = \sum_i m_i r_i^2} \quad \text{and} \quad \boxed{K = \frac{1}{2} I \omega^2}.$$

For continuous bodies,

$$\boxed{I = \int r^2 dm}.$$

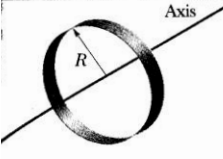
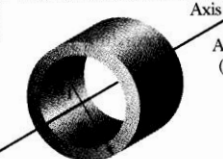
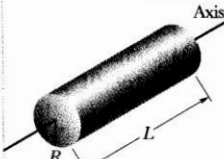
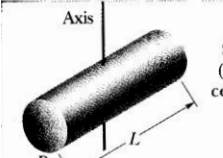
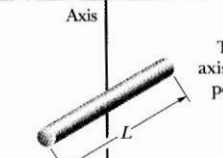
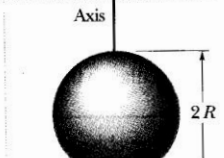
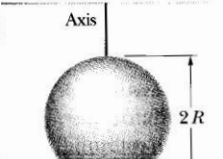
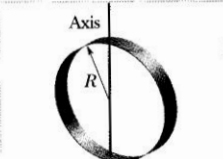

Ring  $I = MR^2$  (axis)  $I = \frac{1}{2}MR^2$  (diameter)

Cylinder  $I = \frac{1}{2}MR^2$

Rod  $I = \frac{1}{12}ML^2$  (centre)  $I = \frac{1}{3}ML^2$  (end)

Sphere  $I = \frac{2}{5}MR^2$

### Some Rotational Inertias

 <p>Hoop about central axis</p> <p><math>I = MR^2</math></p> <p>(a)</p>	 <p>Annular cylinder (or ring) about central axis</p> <p><math>I = \frac{1}{2}M(R_1^2 + R_2^2)</math></p> <p>(b)</p>	 <p>Solid cylinder (or disk) about central axis</p> <p><math>I = \frac{1}{2}MR^2</math></p> <p>(c)</p>
 <p>Solid cylinder (or disk) about central diameter</p> <p><math>I = \frac{1}{4}MR^2 + \frac{1}{12}ML^2</math></p> <p>(d)</p>	 <p>Thin rod about axis through center perpendicular to length</p> <p><math>I = \frac{1}{12}ML^2</math></p> <p>(e)</p>	 <p>Solid sphere about any diameter</p> <p><math>I = \frac{2}{5}MR^2</math></p> <p>(f)</p>
 <p>Thin spherical shell about any diameter</p> <p><math>I = \frac{2}{3}MR^2</math></p> <p>(g)</p>	 <p>Hoop about any diameter</p> <p><math>I = \frac{1}{2}MR^2</math></p> <p>(h)</p>	 <p>Slab about perpendicular axis through center</p> <p><math>I = \frac{1}{12}M(a^2 + b^2)</math></p> <p>(i)</p>

## Parallel Axis Theorem

$$I = I_{\text{cm}} + Mh^2.$$

The rotational inertia of a body about any axis is equal to the rotational inertia ( $= Mh^2$ ) it would have about that axis if all its mass were concentrated at its centre of mass, plus its rotational inertia ( $= I_{\text{cm}}$ ) about a parallel axis through its centre of mass.

### Proof

$$I = \int r^2 dm = \int [(x-a)^2 + (y-b)^2] dm,$$

which can be written as

$$I = \int (x^2 + y^2) dm - 2a \int x dm - 2b \int y dm + \int (a^2 + b^2) dm.$$

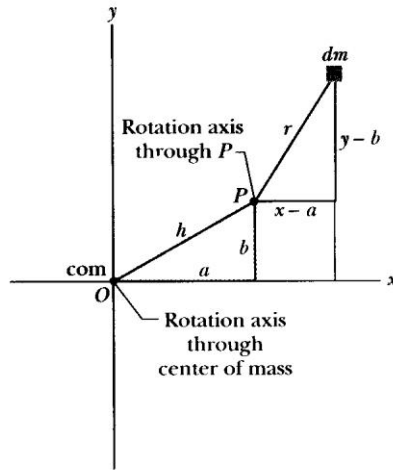


FIG. 10-12 A rigid body in cross section, with its center of mass at  $O$ . The parallel-axis theorem (Eq. 10-36) relates the rotational inertia of the body about an axis through  $O$  to that about a parallel axis through a point such as  $P$ , a distance  $h$  from the body's center of mass. Both axes are perpendicular to the plane of the figure.

In the first term,  $x^2 + y^2 = R^2$ . Hence the first term becomes

$$\int (x^2 + y^2) dm = \int R^2 dm = I_{cm}.$$

In the second and third terms, the position of the centre of mass gives

$$x_{cm} = \frac{1}{M} \int x dm = 0 \quad \text{and} \quad y_{cm} = \frac{1}{M} \int y dm = 0.$$

Hence these terms vanish.

In the last term,  $a^2 + b^2 = h^2$ . Hence the last term becomes

$$\int (a^2 + b^2) dm = \int h^2 dm = Mh^2.$$



## Examples

**10-6** A rigid body consists of two particles of mass  $m$  connected by a rod of length  $L$  and negligible mass.

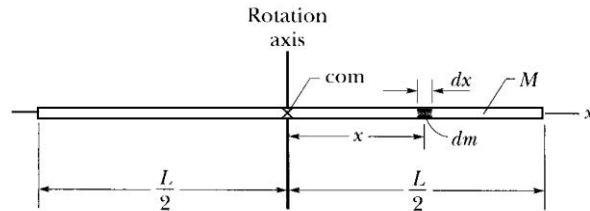
(a) What is the rotational inertia  $I_{\text{cm}}$  about an axis through the center of mass perpendicular to the rod?

(b) What is the rotational inertia  $I$  of the body about an axis through the left end of the rod and parallel to the first axis?

**10-7** Consider a thin, uniform rod of mass  $M$  and length  $L$ .  
 (a) What is the rotational inertia about an axis perpendicular to the rod, through its center of mass?  
 (b) What is the rotational inertia of the rod about an axis perpendicular to the rod through one end?

$$(a) \quad I = \int r^2 dm$$

$$dm = \frac{m}{L} dx$$



**FIG. 10-14** A uniform rod of length  $L$  and mass  $M$ . An element of mass  $dm$  and length  $dx$  is represented.

$$I = \int_{-\frac{L}{2}}^{\frac{L}{2}} x^2 \left( \frac{M}{L} dx \right) = \frac{M}{3L} \left[ x^3 \right]_{-\frac{L}{2}}^{\frac{L}{2}} = \frac{M}{3L} \left[ \left( \frac{L}{2} \right)^3 - \left( -\frac{L}{2} \right)^3 \right] = \frac{ML^2}{12}$$

(b) Using the parallel axis theorem,

$$I = I_{\text{cm}} + Mh^2 = \frac{1}{12} ML^2 + M \left( \frac{L}{2} \right)^2 = \frac{1}{3} ML^2 \quad (\text{ans})$$

## Torque

The ability of  $\vec{F}$  to rotate the body depends on:

- (1) the magnitude of the tangential component  $F_t = F \sin \phi$ ,
- (2) the distance between the point of application and the axis of rotation.

Define the **torque** as

$$\tau = rF \sin \phi.$$

It can be considered as either  $rF_{\perp}$  or  $r_{\perp}F$ . Terms:

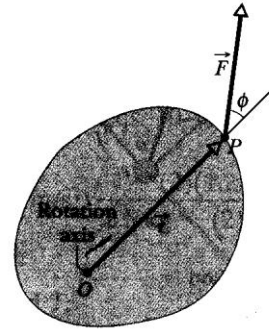
line of action  
moment arm

$\tau$  is positive if it tends to rotate the body counterclockwise.

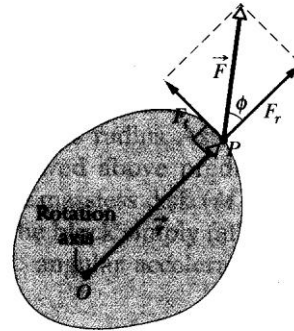
It is negative if it tends to rotate the body clockwise.

Considering the vector direction,

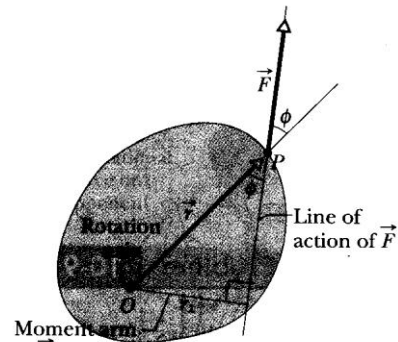
$$\vec{\tau} = \vec{r} \times \vec{F}.$$



(a)



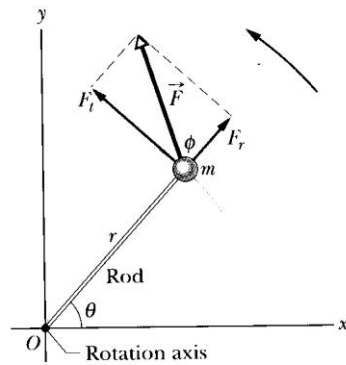
(b)



(c)

**FIG. 10-16** (a) A force  $\vec{F}$  acts at point  $P$  on a rigid body that is free to rotate about an axis through  $O$ ; the axis is perpendicular to the plane of the cross section shown here. (b) The torque due to this force is  $(r)(F \sin \phi)$ . We can also write it as  $rF_t$ , where  $F_t$  is the tangential component of  $\vec{F}$ . (c) The torque can also be written as  $r_{\perp}F$ , where  $r_{\perp}$  is the moment arm of  $\vec{F}$ .

## Newton's Second Law for Rotation



**FIG. 10-17** A simple rigid body, free to rotate about an axis through  $O$ , consists of a particle of mass  $m$  fastened to the end of a rod of length  $r$  and negligible mass. An applied force  $\vec{F}$  causes the body to rotate.

Newton's second law:

$$F_t = ma_t.$$

Torque:

$$\tau = F_t r = ma_t r.$$

Since  $a_t = \alpha r$ , we obtain

$$\tau = m(\alpha r)r = (mr^2)\alpha.$$

Conclusion:

$$\tau = I\alpha.$$

If there are more than one forces,

$$\boxed{\sum \tau = I\alpha.}$$

## Examples

**10-9** A uniform disk of mass  $M = 2.5$  kg and radius  $R = 20$  cm is mounted on a fixed horizontal axle. A block whose mass  $m$  is 1.2 kg hangs from a massless cord that is wrapped around the rim of the disk. Find the acceleration of the falling block, the angular acceleration of the disk, and the tension in the cord. The cord does not slip, and there is no friction at the axle.

Newton's law for the hanging block  
(upward is positive):

$$T - mg = ma \quad (1)$$

Newton's law for the rotating disk  
(anticlockwise is positive):

$$-TR = \frac{1}{2}MR^2\alpha \quad (2)$$

Since  $a = R\alpha$ ,

$$(2): T = -\frac{1}{2}Ma$$

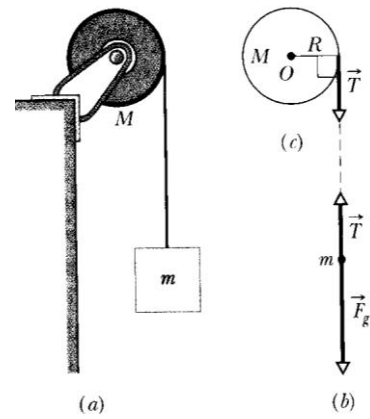
$$(1): -\frac{1}{2}Ma - mg = ma$$

$$-mg = \left(m + \frac{M}{2}\right)a$$

$$a = -\frac{2mg}{M + 2m} = -\frac{(2)(1.2)(9.8)}{2.5 + (2)(1.2)} = -4.8 \text{ ms}^{-2} \quad (\text{ans})$$

$$\alpha = \frac{a}{R} = -\frac{4.8}{0.2} = -24 \text{ rad s}^{-2} \quad (\text{ans})$$

$$T = -\frac{1}{2}Ma = -\left(\frac{1}{2}\right)(2.5)(-4.8) = 6 \text{ N} \quad (\text{ans})$$



**10-10** (Physics of judo) To throw an 80-kg opponent with a basic judo hip throw, you intend to pull his uniform with a force  $\vec{F}$  and a moment arm  $d_1 = 0.30$  m from a pivot point (rotation axis) on your right hip, about which you wish to rotate him with an angular acceleration of  $-6.0 \text{ rad s}^{-2}$ , that is, with a clockwise acceleration. Assume that his rotational inertia  $I$  is  $15 \text{ kg m}^2$ .

- (a) What must the magnitude of  $\vec{F}$  be if you initially bend your opponent forward to bring his centre of mass to your hip?
- (b) What must the magnitude of  $\vec{F}$  be if he remains upright and his weight  $m\vec{g}$  has a moment arm  $d_2 = 0.12$  m from the pivot point?

See Youtube “Judo hip throw”.

(a) Newton’s law for the rotating opponent (anticlockwise is positive):

$$\tau = -d_1 F = I\alpha$$

$$F = -\frac{I\alpha}{d_1} = -\frac{(15)(-6)}{0.3}$$

$$= 300 \text{ N} \quad (\text{ans})$$

(b)

$$\sum \tau = -d_1 F + d_2 mg = I\alpha$$

$$F = -\frac{I\alpha}{d_1} + \frac{d_2 mg}{d_1} = -\frac{(15)(-6)}{0.3} + \frac{(0.12)(80)(9.8)}{0.3} \approx 614 \text{ N}$$

(ans)

Remark: In the correct execution of the hip throw, you should bend your opponent to bring his center of mass to your hip.

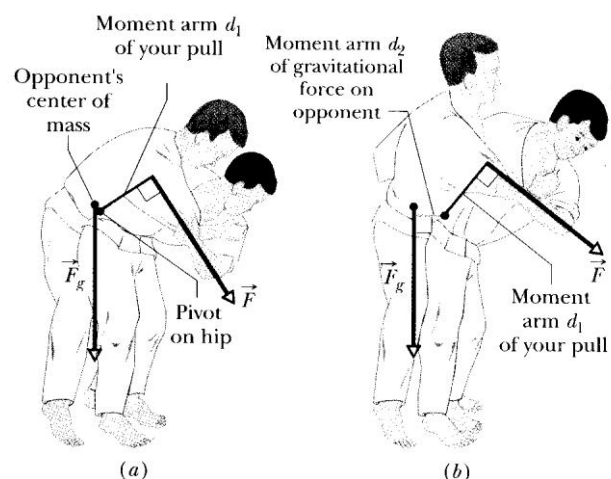


FIG. 10-19 A judo hip throw (a) correctly executed and (b) incorrectly executed.

## Work and Rotational Kinetic Energy

Work done by the force:

$$dW = \vec{F} \cdot d\vec{s} = F_t ds = F_t r d\theta = \tau d\theta.$$

Total work done:

$$W = \int_{\theta_i}^{\theta_f} \tau d\theta.$$

Work-kinetic energy theorem:

$$\tau = I\alpha = I \frac{d\omega}{dt} = I \frac{d\omega}{d\theta} \frac{d\theta}{dt} = I\omega \frac{d\omega}{d\theta} = \frac{d}{d\theta} \left( \frac{1}{2} I\omega^2 \right).$$

Integrating over the angular displacement,

$$W = \int_{\theta_i}^{\theta_f} \tau d\theta = \int_{\theta_i}^{\theta_f} \frac{d}{d\theta} \left( \frac{1}{2} I\omega^2 \right) d\theta = \frac{1}{2} I\omega_f^2 - \frac{1}{2} I\omega_i^2 = \Delta K$$

$$W = \Delta K.$$

## Power

$$P = \frac{dW}{dt} = \tau \frac{d\theta}{dt} = \tau\omega.$$

Some Corresponding Relations for Translational and Rotational Motion

Pure Translation (Fixed Direction)		Pure Rotation (Fixed Axis)	
Position	$x$	Angular position	$\theta$
Velocity	$v = dx/dt$	Angular velocity	$\omega = d\theta/dt$
Acceleration	$a = dv/dt$	Angular acceleration	$\alpha = d\omega/dt$
Mass	$m$	Rotational inertia	$I$
Newton's second law	$F_{\text{net}} = ma$	Newton's second law	$\tau_{\text{net}} = I\alpha$
Work	$W = \int F dx$	Work	$W = \int \tau d\theta$
Kinetic energy	$K = \frac{1}{2} mv^2$	Kinetic energy	$K = \frac{1}{2} I\omega^2$
Power (constant force)	$P = Fv$	Power (constant torque)	$P = \tau\omega$
Work-kinetic energy theorem	$W = \Delta K$	Work-kinetic energy theorem	$W = \Delta K$

## Example

**10-11** As in Example 10-9, a uniform disk of mass  $M = 2.5$  kg and radius  $R = 20$  cm is mounted on a fixed horizontal axle. A block whose mass  $m$  is 1.2 kg hangs from a massless cord that is wrapped around the rim of the disk. What is the rotational kinetic energy  $K$  at  $t = 2.5$  s?



**10-12** A tall, cylindrical chimney will fall over when its base is ruptured. Treat the chimney as a thin rod of length  $L = 55$  m. At the instant it makes an angle of  $\theta = 35^\circ$  with the vertical, what is its angular speed  $\omega_f$ ?

**FIG. 10-20** (a) A cylindrical chimney. (b) The height of its center of mass is determined with the right triangle.

