Physics Enhancement Programme Selection Test 2

Solution

 Construct spherical Gaussian surface with radius *r* and centered at the origin. By symmetry, the field should be radial and with constant magnitude on the surface. Hence Gauss's law implies

$$E \times 4\pi r^{2} = \frac{Q_{\text{enc}}}{\varepsilon_{0}} \Rightarrow \mathbf{E} = E\hat{\mathbf{r}} = \frac{Q_{\text{enc}}}{4\pi\varepsilon_{0}r^{2}}\hat{\mathbf{r}}$$

For $R \ge r \ge 0$, $Q_{\text{enc}} = \rho \times \frac{4}{3}\pi r^{3} \Rightarrow \mathbf{E} = \frac{\rho \times 4\pi r^{3}/3}{4\pi\varepsilon_{0}r^{2}}\hat{\mathbf{r}} = \frac{\rho r}{3\varepsilon_{0}}\hat{\mathbf{r}} = \frac{\rho}{3\varepsilon_{0}}\mathbf{r}$.
For $a > r > R$, $Q_{\text{enc}} = \rho \times \frac{4}{3}\pi R^{3} \Rightarrow \mathbf{E} = \frac{\rho \times 4\pi R^{3}/3}{4\pi\varepsilon_{0}r^{2}}\hat{\mathbf{r}} = \frac{\rho R^{3}}{3\varepsilon_{0}}\hat{\mathbf{r}}$.
For $b > r > a$, $\mathbf{E} = \mathbf{0}$, since it is inside a conductor. (2)

The volume charge density is zero because there can be no volume charge inside a conductor in electrostatics. (0.5)

Construct spherical Gaussian surface centered at the origin and with radius r, where b > r > a. Since the E field is zero inside the conductor, by Gauss's law, the total charge enclosed by this surface should be zero. Hence the total amount of charge on the inner surface is $-\frac{4}{3}\pi\rho R^3$. By symmetry, the charges should be uniformly distributed over the surface. Hence $\sigma = \frac{-4\pi\rho R^3/3}{4\pi a^2} = -\frac{\rho R^3}{3a^2}$. (2.5)

Since the conductor is isolated and carries no net charge, the total amount of charge on the outer surface must be $\frac{4}{3}\pi\rho R^3$. Again, by symmetry, the distribution is uniform. Hence $\sigma = \frac{4\pi\rho R^3/3}{4\pi b^2} = \frac{\rho R^3}{3b^2}$. (1)

Construct spherical Gaussian surface centered at the origin and with radius r, where r > b. Since $Q_{enc} = \frac{4}{3}\pi\rho R^3$, Gauss's law implies

$$\mathbf{E} = \frac{\rho \times 4\pi R^3 / 3}{4\pi\varepsilon_0 r^2} \hat{\mathbf{r}} = \frac{\rho}{3\varepsilon_0} \frac{R^3}{r^2} \hat{\mathbf{r}} .$$
(1.5)

By definition, $V(\mathbf{r}) = -\int_{\infty}^{\mathbf{r}} \mathbf{E} \cdot d\mathbf{l} = -\int_{\infty}^{r} E dr$. For r > b, $V(r) = -\int_{\infty}^{r} \frac{\rho}{3\varepsilon_{0}} \frac{R^{3}}{r^{2}} dr = \frac{\rho R^{3}}{3\varepsilon_{0} r}$. For $b \ge r \ge a$, $V(r) = \frac{\rho R^{3}}{3\varepsilon_{0} b}$. For a > r > R, $V(r) = \frac{\rho R^{3}}{3\varepsilon_{0} b} - \int_{a}^{r} \frac{\rho}{3\varepsilon_{0}} \frac{R^{3}}{r^{2}} dr = \frac{\rho R^{3}}{3\varepsilon_{0} b} - \frac{\rho R^{3}}{3\varepsilon_{0} a} + \frac{\rho R^{3}}{3\varepsilon_{0} r} = \frac{\rho R^{3}}{3\varepsilon_{0}} \left(\frac{1}{b} - \frac{1}{a} + \frac{1}{r}\right)$. For $R \ge r \ge 0$, $V(r) = \frac{\rho R^{3}}{3\varepsilon_{0}} \left(\frac{1}{b} - \frac{1}{a} + \frac{1}{R}\right) - \int_{R}^{r} \frac{\rho r}{3\varepsilon_{0}} dr = \frac{\rho R^{3}}{3\varepsilon_{0}} \left(\frac{1}{b} - \frac{1}{a} + \frac{1}{R}\right) + \frac{\rho}{6\varepsilon_{0}} \left(R^{2} - r^{2}\right)$ $= \frac{\rho}{6\varepsilon_{0}} \left(\frac{2R^{3}}{b} - \frac{2R^{3}}{a} + 3R^{2} - r^{2}\right)$ (4.5)

2. Let the solution of the following configuration be $\Phi_1(x, y)$.

$$A = V = 0$$

$$V = V_0$$

$$V = 0$$

$$V = 0$$

$$X$$

Let the solution of the following configuration be $\Phi_2(x, y)$.

$$A = V = 0$$

$$V = 0$$

$$V = 0$$

$$V = 0$$

$$V = V_0$$

$$A = X$$

Let the solution of the following configuration be $\Phi_3(x, y)$.



Let the solution of the following configuration be $\Phi_4(x, y)$.



Then the solution of the following configuration is

$$\Phi(x, y) = \Phi_{1}(x, y) + \Phi_{2}(x, y) + \Phi_{3}(x, y) + \Phi_{4}(x, y)$$

$$a \qquad V = V_{0}$$

$$V = V_{0}$$

$$V = V_{0}$$

$$V = V_{0}$$

$$x$$

But it is trivial that $\Phi(x, y) = V_0$.

Besides, by symmetry we have

$$\Phi_1(a/2, a/2) = \Phi_2(a/2, a/2) = \Phi_3(a/2, a/2) = \Phi_4(a/2, a/2)$$

Hence

$$\Phi_1(a/2, a/2) = \frac{V_0}{4}$$
 (8)

3. (a)

$$\nabla \times \mathbf{E} = -\frac{\partial \mathbf{B}}{\partial t} \Rightarrow \int_{S} (\nabla \times \mathbf{E}) \cdot d\mathbf{a} = -\int_{S} \frac{\partial \mathbf{B}}{\partial t} \cdot d\mathbf{a}$$
$$\Rightarrow \oint_{C} \mathbf{E} \cdot d\mathbf{l} = -\frac{d}{dt} \int_{S} \mathbf{B} \cdot d\mathbf{a}$$
$$\Rightarrow \mathcal{E} = -\frac{d\Phi}{dt}$$
For $t < 0, \ \Phi = 0, \ \mathcal{E} = -\frac{d\Phi}{dt} = 0.$ For $a/v > t > 0, \ \Phi = Bavt, \ \mathcal{E} = -\frac{d\Phi}{dt} = -vBa.$ For $2a/v > t > a/v, \ \Phi = Ba^{2} - Bavt, \ \mathcal{E} = -\frac{d\Phi}{dt} = vBa.$ For $t > 2a/v, \ \Phi = 0, \ \mathcal{E} = -\frac{d\Phi}{dt} = 0.$

(b)

$$\nabla \cdot \left(\mu_0 \mathbf{J} + \mu_0 \varepsilon_0 \frac{\partial \mathbf{E}}{\partial t} \right) = \mu_0 \left(\nabla \cdot \mathbf{J} + \varepsilon_0 \frac{\partial}{\partial t} \nabla \cdot \mathbf{E} \right)$$
$$= \mu_0 \left(\nabla \cdot \mathbf{J} + \varepsilon_0 \frac{\partial}{\partial t} \frac{\rho}{\varepsilon_0} \right) = \mu_0 \left(\nabla \cdot \mathbf{J} + \frac{\partial \rho}{\partial t} \right) = 0$$
(2)

4. (a) Construct circular Amperian loop with radius *s*. The loop is centered at the wire and lies on a plane perpendicular to it.By symmetry, the H field should be along the tangential direction and with constant magnitude along the wire. Hence, by Ampere's law

$$H \times 2\pi s = I_{f} \Rightarrow \mathbf{H} = \frac{I_{f}}{2\pi s} \hat{\mathbf{\phi}} .$$

For $R > s > 0$, $\mathbf{H} = \frac{\mathbf{B}}{\mu} \Rightarrow \mathbf{B} = \frac{\mu I_{f}}{2\pi s} \hat{\mathbf{\phi}} .$
For $s > R$, $\mathbf{H} = \frac{\mathbf{B}}{\mu_{0}} \Rightarrow \mathbf{B} = \frac{\mu_{0} I_{f}}{2\pi s} \hat{\mathbf{\phi}} .$ (4)

(b) Since $\mathbf{J}_b = \left(\frac{\mu}{\mu_0} - 1\right) \mathbf{J}_f$, inside the cylinder, \mathbf{J}_b is everywhere zero except at

the location of the wire, at which
$$\mathbf{I}_b = \left(\frac{\mu}{\mu_0} - 1\right)\mathbf{I}_f = \left(\frac{\mu}{\mu_0} - 1\right)I_f \hat{\mathbf{z}}$$
. (2)

$$\mathbf{M} = \left(\frac{\mu}{\mu_0} - 1\right) \mathbf{H} = \left(\frac{\mu}{\mu_0} - 1\right) \frac{I_f}{2\pi s} \hat{\mathbf{\phi}} \,. \tag{2}$$

$$\mathbf{K}_{b} = \mathbf{M}(R) \times \hat{\mathbf{s}} = \left(\frac{\mu}{\mu_{0}} - 1\right) \frac{I_{f}}{2\pi R} \hat{\mathbf{\phi}} \times \hat{\mathbf{s}} = -\left(\frac{\mu}{\mu_{0}} - 1\right) \frac{I_{f}}{2\pi R} \hat{\mathbf{z}}.$$
 (2)

5. Let the dipole be

$$\mathbf{p} = p_x \hat{\mathbf{x}} + p_y \hat{\mathbf{y}} + p_z \hat{\mathbf{z}}$$

then the image dipole is

located at
$$(0, 0, -a)$$
.

 $\mathbf{p}' = -p_x \mathbf{\hat{x}} - p_y \mathbf{\hat{y}} + p_z \mathbf{\hat{z}}$

<u>(2)</u>

The total electric field just above the conducting plate is

$$\mathbf{E}(x, y, 0) = \frac{1}{4\pi\varepsilon_0} \frac{1}{\left(x^2 + y^2 + a^2\right)^{3/2}} \\ \left[\frac{3}{x^2 + y^2 + a^2} \left(p_x x + p_y y - p_z a\right) \left(x\hat{\mathbf{x}} + y\hat{\mathbf{y}} - a\hat{\mathbf{z}}\right) - p_x \hat{\mathbf{x}} - p_y \hat{\mathbf{y}} - p_z \hat{\mathbf{z}} \\ + \frac{3}{x^2 + y^2 + a^2} \left(-p_x x - p_y y + p_z a\right) \left(x\hat{\mathbf{x}} + y\hat{\mathbf{y}} + a\hat{\mathbf{z}}\right) + p_x \hat{\mathbf{x}} + p_y \hat{\mathbf{y}} - p_z \hat{\mathbf{z}} \\ = \frac{1}{2\pi\varepsilon_0} \frac{1}{\left(x^2 + y^2 + a^2\right)^{3/2}} \left[\frac{3a}{x^2 + y^2 + a^2} \left(p_z a - p_x x - p_y y\right) - p_z\right] \hat{\mathbf{z}}$$
(3)

The surface charge density can be obtained by

$$\sigma = \varepsilon_0 E = \frac{1}{2\pi} \frac{1}{\left(x^2 + y^2 + a^2\right)^{3/2}} \left[\frac{3a}{x^2 + y^2 + a^2} \left(p_z a - p_x x - p_y y\right) - p_z \right].$$
(2)

6. (a)
$$\mathbf{H}_{s}^{\prime\prime}(a) = \mathbf{H}_{l}^{\prime\prime}(a)$$
. (1)

(b)
$$B_s^{\perp}(a) = B_l^{\perp}(a) \implies \mu_0 \left(H_s^{\perp}(a) + M_0 \cos \theta \right) = \mu H_l^{\perp}(a).$$
 (3)

(c)
$$\mathbf{H}_{l}^{\prime\prime}(b) = \mathbf{H}_{v}^{\prime\prime}(b)$$
. (1)

(d)
$$B_l^{\perp}(b) = B_v^{\perp}(b) \implies \mu H_l^{\perp}(b) = \mu_0 H_v^{\perp}(b).$$
 (2)

7.



(a)
$$R_{23} = \frac{R_2 R_3}{R_2 + R_3} = \frac{(20)(30)}{20 + 30} = 12 \ \Omega.$$
 [2 points]



$$\begin{split} +E &-i_1R_1-i_1R_{23}-i_1R_4=0\\ (12) &-i_1(20)-i_1(12)-i_1(8)=0\\ &i_1=\frac{12}{40}=0.30~\text{A}. \end{split}$$







$$V_{23} = i_1 R_{23} = (0.3)(12) = 3.6 \text{ V.}$$

 $i_2 = \frac{V_2}{R_2} = \frac{3.6}{20} = 0.18 \text{ A.}$ [3 points]

(b)



$$i_{3} = i_{1} - i_{2} = 0.30 - 0.18 = 0.12$$
 A. [1 point]

8.



(a)

$$\begin{split} X_{c} &= \frac{1}{2\pi f_{d}C} = \frac{1}{2\pi (60.0)(15.0 \times 10^{-6})} = 177 \ \Omega. \end{split} \qquad [1.5 \text{ points}] \\ X_{L} &= 2\pi f_{d}L = 2\pi (60.0)(230 \times 10^{-3}) = 86.7 \ \Omega. \end{aligned} \qquad [1.5 \text{ points}] \\ Z &= \sqrt{R^{2} + (X_{L} - X_{C})^{2}} \\ &= \sqrt{(200)^{2} + (86.7 - 177)^{2}} \\ &= 219 \ \Omega. \end{aligned}$$
$$I &= \frac{E_{m}}{Z} = \frac{36.0}{219} = 0.164 \ \text{A}. \end{aligned}$$

(c)

$$\begin{split} \phi &= \tan^{-1} \frac{X_L - X_C}{R} \\ &= \tan^{-1} \frac{86.7 - 177}{200} \\ &= -24.3^\circ = -0.424 \text{ rad.} \end{split}$$

[2 points; deduct 1 point without "-" sign]

9.

(a)

(b)

$$\begin{pmatrix} \text{time interval} \\ \text{relative to you} \end{pmatrix} = \frac{\text{distance relative to you}}{c}$$

$$= \frac{L_0 / \gamma}{c}$$

$$= \frac{9.00 \times 10^{16} / 22.4}{299,792,458}$$

$$= 1.340 \times 10^7 \text{ s} = 0.425 \text{ y}.$$

$$\begin{bmatrix} 4.5 \text{ points} \end{bmatrix}$$

10.

$$\frac{mv^2}{2} = \frac{3kT}{2} \bowtie p = mv = \sqrt{3mkT} = \frac{h}{/}$$
 (2 points)
$$\bowtie / = 1.45 \ 10^{-10}m$$
 (1 points)

Since the wavelength of the thermal neutrons is comparable with the atom separation, they can be diffracted by the crystal. (2 point)

(a)
$$/' - / = \frac{h}{mc} (1 - \cos f) \triangleright / ' - / = 2.4 \ 10^{-12} m$$
 (2 points)
 $\triangleright / ' = 1.24 \ 10^{-11} m$ (1 points)
(b) $DE = hf' - hf$ (1 point)
 $= hc \frac{a}{c} \frac{1}{c} - \frac{1}{c} \frac{\ddot{0}}{g} = -3.88 \ 10^{-15} J$ (1 points)

(c) the kinetic energy of the recoiling electron = $3.88 \cdot 10^{-15} J$ (1 points)

(d)
$$\tan \mathcal{J} = \frac{p}{p'} = \frac{f'}{f} \vartriangleright \mathcal{J} = 51.2^{\circ}$$
 (3 points)

The electron make an angle $90 - 51.2 = 38.8^{\circ}$ with the x-axis. (1 point)

12.

$$\frac{1}{s_1} + \frac{1}{s_1'} = \frac{1}{f} \vartriangleright \frac{1}{s_1'} = \frac{1}{10} - \frac{1}{15} \vartriangleright s_1' = 30cm \quad (1 \text{ point})$$

(i.e. the image is 30cm to the right of the first lens

or it is 10cm to the right of the second lens)

$$s_{2} = -10cm$$
 (1 point)
$$\frac{1}{s_{2}} + \frac{1}{s_{2}} = \frac{1}{f} \vartriangleright \frac{1}{s_{2}} = \frac{1}{10} + \frac{1}{10} \bowtie s_{2} = 5cm$$
 (1 point)

the image is 5cm to the right of the second lens which is real (1 point).

The magnification m = $\overset{\mathfrak{R}}{\underset{0}{\overset{\circ}{\leftrightarrow}}} - \frac{s_1'\overset{0}{\underset{1}{\overset{\circ}{\otimes}}}}{s_1}\overset{0}{\underset{0}{\overset{\circ}{\otimes}}} - \frac{s_2'\overset{0}{\underset{2}{\overset{\circ}{\otimes}}}}{s_2}\overset{0}{\underset{0}{\overset{\circ}{\otimes}}} = -1.5$ (the image is magnified and inverted) (1 point)

13.

$$\mathcal{J}_{0} = 30^{\circ} \quad (1 \text{ point})
\bowtie n \sin \mathcal{J}_{1} = \sin \mathcal{J}_{0} \bowtie \mathcal{J}_{1} = 19.2 \quad (1 \text{ point})
60 + (90 - \mathcal{J}_{1}) + (90 - \mathcal{J}_{2}) = 180 \bowtie \mathcal{J}_{2} = 60 - \mathcal{J}_{1} = 40.8 \quad (2 \text{ point})
\sin \mathcal{J}_{3} = n \sin \mathcal{J}_{2} \bowtie \mathcal{J}_{3} = 83^{\circ} \quad (1 \text{ point})$$

11.