

**Pan Pearl River Delta Physics Olympiad 2012**  
**2012 年泛珠三角及中华名校物理奥林匹克邀请赛**  
**Part-1 (Total 6 Problems) 卷-1 (共 6 题)**  
 (9:00 am – 12:00 pm, 02-02-2012)

**Q1 (5 points)**

$$E = \gamma mc^2, \text{ where } \gamma \equiv \frac{1}{\sqrt{(1-\beta)(1+\beta)}} \approx \frac{1}{\sqrt{2(1-\beta)}}.$$

$$\text{So } mc^2 = E\sqrt{2(1-\beta)} < 10 \cdot \sqrt{2 \cdot 2 \cdot 10^{-9}} = 632 \text{ eV. (4 points)}$$

Upper limit 上限. (1 point)

**Q2 (10 points)**

设 X-Y 平面为水平面，单摆平衡时质点在 (0, 0)，质点现在位置是 (x, y)。利用小幅振动近似的结果，我们知道细绳的张力等于质点的重力，张力在 X-Y 平面的投影的大小为  $mg \frac{\sqrt{x^2 + y^2}}{L}$ ，方向指向 (0, 0)。因此，X 方向的分力为

$$F_x = -mg \frac{\sqrt{x^2 + y^2}}{L} \frac{x}{\sqrt{x^2 + y^2}} = -\frac{x}{L} mg \quad (\text{i}).$$

$$\text{Y 方向的分力为 } F_y = -mg \frac{\sqrt{x^2 + y^2}}{L} \frac{y}{\sqrt{x^2 + y^2}} = -\frac{y}{L} mg \quad (\text{ii}).$$

$$\text{动力学方程为 } -\frac{x}{L} g = \ddot{x}, \quad -\frac{y}{L} g = \ddot{y}.$$

通解为

$$x(t) = A_x \cos(\omega t) + B_x \sin(\omega t), \quad y(t) = A_y \cos(\omega t) + B_y \sin(\omega t),$$

$$\vec{r}(t) = x(t)\vec{x}_0 + y(t)\vec{y}_0. \quad \omega = \sqrt{g/L}. \text{ (2 points)}$$

设一般的初始条件  $\vec{r}(0) = X_0\vec{x}_0 + Y_0\vec{y}_0$ ,  $\dot{\vec{r}}(0) = v_{x0}\vec{x}_0 + v_{y0}\vec{y}_0$ , 则

$$x(t) = X_0 \cos(\omega t) + \frac{v_{x0}}{\omega} \sin(\omega t), \quad y(t) = Y_0 \cos(\omega t) + \frac{v_{y0}}{\omega} \sin(\omega t). \text{ (2 points)}$$

As we only ask for examples, there can be many different ways. 题目只要求给出例子，所以可以有多种答案。

$$(a) \quad X_0 = D/2, \text{ others are 0 其余为 0. Then 则 } x(t) = \frac{D}{2} \cos(\omega t), \quad y(t) = 0. \text{ (2 points)}$$

$$(b) \quad \vec{r}(0) = R\vec{x}_0, \quad \dot{\vec{r}}(0) = R\omega\vec{y}_0. \text{ Then } x(t) = R \cos(\omega t), \quad y(t) = R \sin(\omega t). \text{ (2 points)}$$

$$(c) \quad \vec{r}(0) = a\vec{x}_0, \quad \dot{\vec{r}}(0) = b\omega\vec{y}_0. \text{ Then } x(t) = a \cos(\omega t), \quad y(t) = b \sin(\omega t). \text{ (2 points)}$$

**Q3 (9 points)**

In the rotating reference frame, the force along the radial direction is 在跟着卫星转的旋转参照系里，沿半径方向的力为：

$$F_r(r) = m\omega^2 r - \frac{GMm}{r^2}, \text{ and } F_r(R) = m\omega_0^2 R - \frac{GMm}{R^2} = 0.$$

After the impact, the total angular momentum is still conserved 碰撞后, 角动量仍然守恒:

$$0 = d(\omega r^2)_{r=R} = R^2 d\omega + 2\omega_0 R dr.$$

Let the orbit radius change by  $dr$ , then 令轨道的变化为  $dr$

$$dF_r(R) = m\omega_0^2 dr + 2\omega_0 R d\omega + m \frac{2GMm}{R^3} dr = -m\omega_0^2 dr.$$

(Note that without taking into account the change of  $\omega$ , the force is positive and the balance is unstable. 若漏了考虑  $\omega$  的变化, 则力的变化是正的, 原来的轨道运动变得不平衡了, 那是不对的。)

This is a SHM with 上式结果显示卫星的径向运动是简谐振动, 力常数为  $k = m\omega_0^2$ , so 因此振动频率为  $\omega = \omega_0$ .

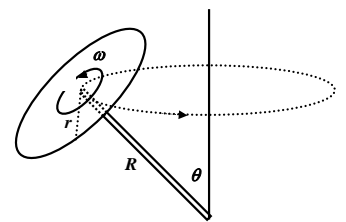
The initial condition is 初始条件为  $v_0 = I / m$ , and  $dr = 0$ . So  $dr(t) = \frac{-I}{m\omega_0} \sin(\omega_0 t)$ .

$$x(t) = [R - \frac{I}{m\omega_0} \sin(\omega_0 t)] \cos(\omega_0 t), \quad y(t) = [R - \frac{I}{m\omega_0} \sin(\omega_0 t)] \sin(\omega_0 t).$$

**Q4 (6 points)**

Angular momentum is 角动量为  $J = mr^2 \omega / 2$ .

Torque 力矩  $\tau = mgR \sin \theta$ , and pointing perpendicular to the paper plane 方向垂直于纸面.



The change of angular momentum is 角动量的变化为  $\Delta J = (J \sin \theta) \Delta \phi$ .

$$\tau = \frac{\Delta J}{\Delta t} = (J \sin \theta) \frac{\Delta \phi}{\Delta t} = (J \sin \theta) \Omega. \text{ So } \Omega = \frac{2gR}{r^2 \omega}.$$

**Q5 (10 points)**

Let the image charge be at  $x$  on the X-axis 放个镜像电荷在 X-轴上  $x$  点处。

$$\Phi(R) = \frac{1}{4\pi\epsilon_0} \left( \frac{q}{\sqrt{R^2 + x^2 + 2Rx \cos \theta}} - \frac{q_1}{\sqrt{R^2 + b^2 + 2Rb \cos \theta}} \right) = 0$$

$$\Rightarrow q^2 (R^2 + b^2 + 2Rb \cos \theta) = q_1^2 (R^2 + x^2 + 2Rx \cos \theta).$$

上式对任何角度  $\theta$  都成立, 所以有

$$q^2 (R^2 + b^2) = q_1^2 (R^2 + x^2) \quad \text{(i),}$$

$$q^2 b = q_1^2 x \quad \text{(ii).}$$

将  $b$  代掉, 得

$$x^2 q_1^4 - q^2 (R^2 + x^2) q_1^2 + q^4 R^2 = 0 \quad \text{(iii).}$$

解上式, 得两个解。第一个是  $q_1 = -q$ ,  $b = x$ , 即把球外的点电荷中和掉。因为这个镜像电荷没有放在球面内, 不在所考虑的解的空间以外, 所以不能用, 舍去。第二个是  $q_1 = -Rq/x$ ,  $b = R^2/x$ 。这

就是镜像电荷的值和位置。(It is OK if answers are given without derivations.若没有上述推导而只有答案也得全分。)

In order to make the potential on the plane zero, we need two more image charges, namely  $q_2 = -q_1$  at  $-R^2/x$ , and  $q_3 = -q$  at  $-3R/2$ . One can see that  $q_2$  and  $q_3$  combined will make the potential on the sphere surface zero. One can see that  $q_2$  and  $q_3$  combined will make the potential on the sphere surface zero. 为了使平面的电势为 0, 我们需要另外两的镜像电荷,  $q_2 = -q_1$  在  $-R^2/x$ , and  $q_3 = -q$  在  $-3R/2$ . 而  $q_2$ 、 $q_3$  合起来也使球面的电势为 0。(2 points)

$$\text{电荷受的力为: } F_1 = \frac{qq_1}{4\pi\epsilon_0(x-b)^2} = \frac{-q^2}{4\pi\epsilon_0} \frac{R}{x(x-R^2/x)^2} = \frac{-q^2}{4\pi\epsilon_0} \frac{xR}{(x^2-R^2)^2}.$$

现在求  $q$  的电势, 也就是把  $q$  从无穷远拉到现在位置所需的能量。用作用力做功的方法,

$$W_1 = \int_{3R/2}^{\infty} F_1 dx = \frac{-q^2 R}{4\pi\epsilon_0} \int_{3R/2}^{\infty} \frac{x}{(x^2-R^2)^2} dx = \frac{-q^2}{8\pi\epsilon_0} \left( \frac{R}{d^2-R^2} \right) = \frac{-q^2}{8\pi\epsilon_0 R} \left( \frac{4}{9-4} \right) = \frac{-q^2}{10\pi\epsilon_0 R}. \quad (4 \text{ points})$$

读者请注意, 此势能和直接用电势所得的值是不同的。 $q$  所在位置的电势由  $q_1$  产生, 其值为

$$U = \frac{q_1}{4\pi\epsilon_0(d-b)} = \frac{-q}{4\pi\epsilon_0(d^2-R^2)} = \frac{-q}{5\pi\epsilon_0 R}. \text{ 所以 } q \text{ 的势能为 } W = \frac{-q^2}{5\pi\epsilon_0 R}.$$

正确答案是错误答案的一半。两者差别的主要原因, 是因为镜像电荷的值随真电荷的位置而变。所以计算电荷势能最可靠的方法是用作用力的路径积分来做。

$$W_2 = \int_{3R/2}^{\infty} F_2 dx = \frac{q^2 R}{4\pi\epsilon_0} \int_{3R/2}^{\infty} \frac{x}{(x^2+R^2)^2} dx = \frac{q^2}{8\pi\epsilon_0 R} \left( \frac{4}{9+4} \right) = \frac{q^2}{26\pi\epsilon_0 R}.$$

$$W_3 = \int_{3R/2}^{\infty} F_3 dx = \int_{3R/2}^{\infty} \frac{q^2}{16\pi\epsilon_0 x^2} dx = -\frac{q^2}{24\pi\epsilon_0 R}. \quad (2 \text{ points})$$

$$W = W_1 + W_2 + W_3 = \frac{q^2}{2\pi\epsilon_0 R} \left( -\frac{1}{5} + \frac{1}{13} - \frac{1}{12} \right) = \frac{-161}{1560} \frac{q^2}{\pi\epsilon_0 R}. \quad (1 \text{ point})$$

### Q6 (10 points)

(a) Let  $c = \frac{3}{2}$ , so that  $c_v = cnR$ . The expansion is an adiabatic process, 膨胀过程为绝热过程, 所以有

$$PV^{1+1/c} = \text{con}.$$

The work done to the piston is 对活塞做的功为

$$W = \int_{V_0}^{\kappa V_0} PdV = P_0 V_0^{1+1/c} \int_{V_0}^{\kappa V_0} V^{-1-1/c} dV = cP_0 V_0 (1 - \kappa^{-1/c}) = \frac{3}{2} nRT_0 (1 - \kappa^{-2/3}). \quad (2 \text{ points})$$

Maximum work is 最大功为  $W_{\max} = \frac{3}{2}nRT_0$ , which is the total internal energy of the gas 等于气体的总内能, 也就是气体可做的最大功. (1 point)

$$(b) \quad \frac{\Delta S}{nR} = \ln(\kappa) + c \ln(T/T_0) = \ln(\kappa) + c \ln\left(\frac{1}{T_0}\left(T_0 - \frac{\eta W}{cnR}\right)\right)$$

Put in the answer in (a) we get 代入(a) 里功的表达式, 得  $\frac{\Delta S}{nR} = \ln(\kappa) + \frac{3}{2} \ln(1 - \eta + \eta\kappa^{-2/3})$ . (5 points)

To see if  $\Delta S$  is positive, let  $\Delta S = cnR \ln(m)$ , so  $m - 1 = (1 - \eta)(\kappa^{-2/3} - 1) \geq 0$ , and  $\Delta S > 0$ . (1 point)

(c) When  $\eta = 1$ ,  $\Delta S = 0$ , which is consistent with the fact that there is no net loss in internal energy of the gas. 当  $\eta = 1$ ,  $\Delta S = 0$ , which is consistent with the fact that there is no net loss in internal energy of the gas 符合气体自由膨胀时熵不变这一结果。 (1 point)

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**Part-2 (Total 3 Problems) 卷-2 (共 3 题)**  
(2:30 pm – 5:30 pm, 02-02-2011)

**Q1 Stella Interferometer (8 points)**

- (a) The difference between the path differences of the two pairs of split waves is  $\Delta = D \cdot \delta\theta$ . When  $\Delta = \lambda/2$ , the bright fringes of one star coincide with the dark fringes of the other. 两个平面波在两个入口镜的光程差的差为  $\Delta = D \cdot \delta\theta$ . 当  $\Delta = \lambda/2$  时, 一颗星的亮条纹与另一颗星的暗条纹刚好重叠, 所以:  $D = \frac{\lambda}{2\delta\theta} = \frac{180 \cdot 5.0 \cdot 10^{-7}}{2 \cdot 3.14 \cdot 3 \cdot 10^{-6}} m = 4.8 m$ . (6 points)

(b)  $\delta\theta = \frac{1.22\lambda}{D}$ , so  $D = \frac{1.22\lambda}{\delta\theta} = \frac{1.22 \cdot 180 \cdot 5.0 \cdot 10^{-7}}{3.14 \cdot 3 \cdot 10^{-6}} m = 12 m$ . (2 points)

**Q2 Y-particle (12 points)**

- (a) The two B-mesons should have the same momentum and energy because the original Y-meson is at rest. The kinetic energy of one B-meson is

$$E_k = \frac{1}{2}(10.58 - 2 \times 5.28) = 0.01 \text{ GeV, which is much less than the rest energy of B-mesons.}$$

So we can use  $E_k = \frac{1}{2} m_B v_0^2$ . Putting the numbers in we get  $v_0 = \sqrt{\frac{0.01}{5.28}} c = 0.0615c$ . As the B-mesons are moving at low speed their lifetime change can be ignored. So

$$L_0 = v_0 \tau_0 = 3 \times 10^8 \times 0.0615 \times 1.5 \times 10^{-12} = 0.028 \text{ mm.}$$

由于 Y 介子是静止的, 所以两个 B 介子的动量相等, 方向相反。B 介子的动能为

$$E_k = \frac{1}{2}(10.58 - 2 \times 5.28) = 0.01 \text{ GeV, 比它的静止质量小得多, 所以可用经典力学}$$

$$E_k = \frac{1}{2} m_B v_0^2 \text{ 来求它的速度。将数值代入后得 } v_0 = \sqrt{\frac{0.01}{5.28}} c = 0.0615c. \text{ B 介子寿命因运动}$$

而延长的效应可忽略, 因此

$$L_0 = v_0 \tau_0 = 3 \times 10^8 \times 0.0615 \times 1.5 \times 10^{-12} = 0.028 \text{ mm. (2 points)}$$

- (b) Rough estimation: the speed of the B-mesons should be  $v_0 L / L_0 = 0.44c$ , so we need precise formula which takes into account the lifetime change. 粗略估计: B 介子的速度为  $v_0 L / L_0 = 0.44c$ , 所以必须考虑相对论效应。

$$L = \frac{v\tau_0}{\sqrt{1-(v/c)^2}}. \text{ So } v = \frac{c}{\sqrt{(c\tau_0/L)^2 + 1}} = 0.406c.$$

$$P = \gamma m_B v = 0.406 \times 1.094 \times 5.28 (\text{GeV}/c) = 2.35 \text{ GeV}/c \text{ (3 points)}$$

- (c) The Y-mesons move with the center-of-mass (CoM) frame of the B-mesons. From (a) we get the speed of the B-mesons in that frame as  $v_0 = 0.0615c$ . From (b) we know that in the laboratory frame the speed of the B-mesons is  $v = 0.406c$ . Let the relative speed

between the CoM frame and the laboratory frame be  $-u$ , which is also the speed of the Y-mesons in the laboratory frame, then  $v = \frac{v_0 + u}{1 + v_0 u / c^2}$ . Solving it we get  $u = 0.336c$ . Y

介子的速度与 B 介子质心的速度一致, 由 (a) 我们知道 B 介子在质心参照系的速度为  $v_0 = 0.0615c$ . 由 (b) 我们知道 B 介子在实验室参照系里的速度为  $v = 0.406c$ . 令质心参照系相对于实验室的速度 (也是 Y 介子在实验室的速度) 为  $-u$ , 则

$v = \frac{v_0 + u}{1 + v_0 u / c^2}$ . 由此得  $u = 0.336c$ . The total energy of the Y-mesons in the

laboratory frame is Y 介子的总能量为  $E_Y = m_Y c^2 / \sqrt{1 - (u/c)^2} = 1.062 m_Y c^2$ . (4 points)

- (d) In the CoM an electrons has equal and opposite momentum as a positron, while together they must have the energy to create a Y-meson. The momentum 4-vector is

then  $\begin{pmatrix} 0 \\ m_Y c \end{pmatrix}$  in CoM. 在 Y 介子参照系里电子和正电子的动量相等, 方向相反, 总

能量等于 Y 介子的静质量。因此 Y 介子的动量-能量 4 矢为  $\begin{pmatrix} 0 \\ m_Y c \end{pmatrix}$ 。In the

laboratory frame the momentum 4-vector is 在实验室参照系, 该 4 矢为

$$\begin{pmatrix} \gamma & \beta\gamma \\ -\beta\gamma & \gamma \end{pmatrix} \begin{pmatrix} 0 \\ m_Y c \end{pmatrix} = \begin{pmatrix} \beta\gamma m_Y c \\ \gamma m_Y c \end{pmatrix}.$$

Let the momenta of electron and positron be  $P_1$  and  $P_2$ , respectively, and also note that as the electron (and positron) rest energy (0.511 MeV) is thousands of times less than that of the Y-meson, so the energy of an electron is simply  $cP_1$ . 令电子和正电子的动量分别为  $P_1$ 、 $P_2$ , 又由于现在电子的能量远大于它的静止质量 (0.511 MeV), 所以它的能量为  $cP_1$ 。We then have 由此我们得

$$\begin{pmatrix} P_1 - P_2 \\ P_1 + P_2 \end{pmatrix} = \begin{pmatrix} \beta\gamma m_Y c \\ \gamma m_Y c \end{pmatrix}, \text{ where 其中 } \beta \equiv u/c.$$

The electron energy in the laboratory frame is then 电子在实验室的能量为

$$cP_1 = \frac{1}{2}(1 + \beta)\gamma m_Y c^2 = \frac{1}{2} \times 1.336 \times 1.062 \times m_Y c^2 = 0.709 m_Y c^2 = 7.51 \text{ GeV}.$$

正电子的能量为

$$cP_2 = \frac{1}{2}(1 - \beta)\gamma m_Y c^2 = \frac{1}{2} \times 0.664 \times 1.062 \times m_Y c^2 = 0.353 m_Y c^2 = 3.73 \text{ GeV, or vice versa.}$$

(3 points)

### Q3 Penning Trap (30 points)

(a)  $qBv = m \frac{v^2}{r_0}$ , so  $\omega_c = \frac{v}{r_0} = \frac{qB}{m}$ .  $E_k = \frac{1}{2}mv^2 = \frac{1}{2}m\omega_c^2 r_0^2$ . (1 point)

(b)  $x = \frac{D}{2} - r \cos(\omega_c t - \phi)$ ,  $Q_1 - Q_2 = -2q \frac{r}{D} \cos(\omega_c t - \phi)$ .

$I = \frac{dQ}{dt} = 2q \frac{r_0 \omega_c}{D} \sin(\omega_c t - \phi)$ . Larger  $r_0$  leads to larger current. (2 points)

- (c) The energy gained in each cycle is 每周期得到的能量为

$$dE_k = \int_{T_c} qE_0 \cos(\omega_c t) \cdot r\omega_c \cos(\omega_c t) dt = \frac{1}{2} qE_0 r\omega_c T_c, \text{ and } \frac{dE_k}{dt} = \frac{1}{2} qE_0 r\omega_c \quad (1 \text{ point})$$

On the other hand, from (a) we have 由(a) 我们得  $\frac{dE_k}{dt} = m\omega_c^2 r \frac{dr}{dt}$ . (1 point)

$$\text{So } m\omega_c^2 r \frac{dr}{dt} = \frac{1}{2} qE_0 r\omega_c, \text{ and } R = r_0 + \frac{qE_0}{2m\omega_c} T \quad (1 \text{ point})$$

(d)  $E_z = -2V_0 \frac{z}{z_0^2}$ .  $\omega_z^2 = \frac{2qV_0}{mz_0^2}$ . (2 points)

(e)  $\nabla^2 V(\vec{r}) = 0$  so  $\beta = -1/2$ . (1 point)

Use  $\omega_c$  and  $\omega_z$  as known for the remaining part of the question.

(f) The electric field in the x-y plane is 沿 X-Y 平面的电场为  $\vec{E} = -V_0 \frac{x\vec{x}_0 + y\vec{y}_0}{z_0^2}$ . (1 point)

The equation is 粒子的运动方程为

$$m(\ddot{x}\vec{x}_0 + \ddot{y}\vec{y}_0) = eV_0 \frac{x\vec{x}_0 + y\vec{y}_0}{z_0^2} + eB(\dot{x}\vec{x}_0 + \dot{y}\vec{y}_0) \times \vec{z}_0 = eV_0 \frac{x\vec{x}_0 + y\vec{y}_0}{z_0^2} - eB(\dot{x}\vec{y}_0 - \dot{y}\vec{x}_0)$$

(1 point)

$$\ddot{x} - \omega_c \dot{y} - \frac{1}{2} \omega_z^2 x = 0$$

$$\ddot{y} + \omega_c \dot{x} - \frac{1}{2} \omega_z^2 y = 0 \quad (2 \text{ points})$$

- (g) Multiply  $i$  to the first equation, and add to the second one, 将第一式乘  $i$  后与第二式相加, 得 we get  $\ddot{u} + ia\dot{u} + bu = 0$ , where  $a = \omega_c$  and  $b = -\frac{1}{2} \omega_z^2$ . (2 points)

(h)  $\omega_{\pm} = \frac{\omega_c \pm \sqrt{\omega_c^2 - \omega_z^2}}{2}$  (2 points)

(i)  $x(0) = R$ ,  $y(0) = 0$ ,  $\dot{x}(0) = 0$ ,  $\dot{y}(0) = -R\omega_c$  (2 points)

Then  $R = A_+ + A_-$ ,  $R\omega_c = \omega_+ A_+ + \omega_- A_-$  (1 point)

$$\text{Solving the equations we get 解上述方程, 得 } A_- = \frac{\omega_+ - \omega_c}{\omega_+ - \omega_-} R = -\frac{1}{2} \left( \frac{\omega_z}{\omega_c} \right)^2 R,$$

$$A_+ = \frac{\omega_c - \omega_-}{\omega_+ - \omega_-} R = R \quad (2 \text{ points})$$

(j)  $\tilde{x} = x \cos(\Omega t) + y \sin(\Omega t)$ ,  $\tilde{y} = y \cos(\Omega t) - x \sin(\Omega t)$ . So  $\tilde{u} = u e^{-i\Omega t}$ . (3 points)

(k)  $\tilde{u} = u e^{i\omega t} = A_- + A_+ e^{-i(\omega_+ - \omega_-)t}$ , which is a circle centered at  $x = A_-$ . 这是一个中心在  $x = A_-$  的圆。 (1 point)

- (l) (k)中的圆心绕原点作圆周运动。 (1 point)

(m)  $\tilde{u} = A_+ e^{-i\omega_1 t} + A_- e^{i\omega_1 t}$  where  $\omega_1 = \sqrt{\omega_c^2 - \omega_z^2}$ . (1 point)

So  $\tilde{x} = (A_+ + A_-) \cos(\omega_1 t)$ ,  $\tilde{y} = (A_+ - A_-) \sin(\omega_1 t)$ , which is an ellipse with  $(A_+ + A_-)$  being one axis and  $(A_+ - A_-)$  being the other. Under more special conditions, the ellipse can become a line or a circle. 这是个椭圆，在特定条件下可变成圆( $A_-$  或  $A_+$  等于 0), 或直线( $A_- = \pm A_+$ ) (2 points)

《THE END 完》