## THEORETICAL PROBLEM No. 1

## EVOLUTION OF THE EARTH-MOON SYSTEM

## SOLUTIONS

## 1. Conservation of Angular Momentum

| 1a | $L_{1}=I_{E} \omega_{E 1}+I_{M 1} \omega_{M 1}$ | 0.2 |
| :---: | :---: | :---: |
| 1b | $L_{2}=I_{E} \omega_{2}+I_{M 2} \omega_{2}$ | 0.2 |
| 1c | $I_{E} \omega_{E 1}+I_{M 1} \omega_{M 1}=I_{M 2} \omega_{2}=L_{1}$ | 0.3 |

2. Final Separation and Angular Frequency of the Earth-Moon System.

| 2 a | $\omega_{2}^{2} D_{2}^{3}=G M_{E}$ | 0.2 |
| :--- | :--- | :--- |


| 2 b | $D_{2}=\frac{L_{1}^{2}}{G M_{E} M_{M}^{2}}$ | 0.5 |
| :--- | :--- | :--- |

$$
\begin{array}{|l|l|l|}
\hline 2 \mathrm{c} & \omega_{2}=\frac{G^{2} M_{E}^{2} M_{M}^{3}}{L_{1}^{3}} & 0.5 \\
\hline
\end{array}
$$

| 2 d | The moment of inertia of the Earth will be the addition of the moment of <br> inertia of a sphere with radius $r_{o}$ and density $\rho_{o}$ and of a sphere with <br> radius $r_{i}$ and density $\rho_{i}-\rho_{o}:$ <br> $I_{E}=\frac{2}{5} \frac{4 \pi}{3}\left[r_{o}^{5} \rho_{o}+r_{i}^{5}\left(\rho_{i}-\rho_{o}\right)\right]$. | 0.5 |
| :--- | :--- | :--- |


| 2 e | $I_{E}=\frac{2}{5} \frac{4 \pi}{3}\left[r_{o}^{5} \rho_{o}+r_{i}^{5}\left(\rho_{i}-\rho_{o}\right)\right]=8.0 \times 10^{37} \mathrm{~kg} \mathrm{~m}^{2}$ | 0.2 |
| :--- | :--- | :--- |


| 2 f | $L_{1}=I_{E} \omega_{E 1}+I_{M 1} \omega_{M 1}=3.4 \times 10^{34} \mathrm{~kg} \mathrm{~m}^{2} \mathrm{~s}^{-1}$ | 0.2 |
| :--- | :--- | :--- |


| 2 g | $D_{2}=5.4 \times 10^{8} \mathrm{~m}$, that is $D_{2}=1.4 D_{1}$ | 0.3 |
| :--- | :--- | :--- |


| 2 h | $\omega_{2}=1.6 \times 10^{-6} \mathrm{~s}^{-1}$, that is, a period of 46 days. | 0.3 |
| :--- | :--- | :--- |


| 2 i | Since $I_{E} \omega_{2}=1.3 \times 10^{32} \mathrm{~kg} \mathrm{~m}^{2} \mathrm{~s}^{-1}$ and $I_{M 2} \omega_{2}=3.4 \times 10^{34} \mathrm{~kg} \mathrm{~m}^{2} \mathrm{~s}^{-1}$, the <br> approximation is justified since the final angular momentum of the Earth <br> is $1 / 260$ of that of the Moon. | 0.2 |
| :--- | :--- | :--- |

3. How much is the Moon receding per year?

| 3 a | Using the law of cosines, the magnitude of the force produced by the mass <br> $m$ closest to the Moon will be: <br> $F_{c}=\frac{G m M_{M}}{D_{1}^{2}+r_{o}^{2}-2 D_{1} r_{o} \cos (\theta)}$ | 0.4 |
| :--- | :--- | :--- |

3b $\quad$ Using the law of cosines, the magnitude of the force produced by the mass 0.4 $m$ farthest to the Moon will be:
$F_{f}=\frac{G m M_{M}}{D_{1}^{2}+r_{o}^{2}+2 D_{1} r_{o} \cos (\theta)}$

| $3 c$ | Using the law of sines, the torque will be <br> $\tau_{c}=F_{c} \frac{\sin (\theta) r_{0} D_{1}}{\left[D_{1}^{2}+r_{o}^{2}-2 D_{1} r_{o} \cos (\theta)\right]^{1 / 2}}=\frac{G m M_{M} \sin (\theta) r_{0} D_{1}}{\left[D_{1}^{2}+r_{o}^{2}-2 D_{1} r_{o} \cos (\theta)\right]^{3 / 2}}$ | 0.4 |
| :--- | :--- | :--- |


| 3 d | $\begin{array}{l}\text { Using the law of sines, the torque will be } \\ \tau_{f}=F_{f} \frac{\sin (\theta) r_{0} D_{1}}{\left[D_{1}^{2}+r_{o}^{2}+2 D_{1} r_{o} \cos (\theta)\right]^{1 / 2}}=\frac{G m M_{M} \sin (\theta) r_{0} D_{1}}{\left[D_{1}^{2}+r_{o}^{2}+2 D_{1} r_{o} \cos (\theta)\right]^{3 / 2}}\end{array}$ | 0.4 |
| :--- | :--- | :--- |

$$
\begin{array}{|l|l|l|}
\hline 3 \mathrm{e} & \tau_{c}-\tau_{f}=G m M_{M} \sin (\theta) r_{0} D_{1}^{-2}\left(1-\frac{3 r_{o}^{2}}{2 D_{1}^{2}}+\frac{3 r_{o} \cos (\theta)}{D_{1}}-1+\frac{3 r_{o}^{2}}{2 D_{1}^{2}}+\frac{3 r_{o} \cos (\theta)}{D_{1}}\right) & 1.0 \\
& =\frac{6 G m M_{M} r_{o}^{2} \sin (\theta) \cos (\theta)}{D_{1}^{3}} & \\
\hline
\end{array}
$$

| 3f | $\tau=\frac{6 G m M_{M} r_{o}^{2} \sin (\theta) \cos (\theta)}{D_{1}^{3}}=4.1 \times 10^{16} \mathrm{~N} \mathrm{~m}$ | 0.5 |
| :--- | :--- | :--- |


| 3 g | Since $\omega_{M 1}^{2} D_{1}^{3}=G M_{E}$, we have that the angular momentum of the Moon is | 1.0 |
| :--- | :--- | :--- |

$I_{M 1} \omega_{M 1}=M_{M} D_{1}^{2}\left[\frac{G M_{E}}{D_{1}^{3}}\right]^{1 / 2}=M_{M}\left[D_{1} G M_{E}\right]^{1 / 2}$
The torque will be:
$\tau=\frac{M_{M}\left[G M_{E}\right]^{1 / 2} \Delta\left(D_{1}^{1 / 2}\right)}{\Delta t}=\frac{M_{M}\left[G M_{E}\right]^{1 / 2} \Delta D_{1}}{2\left[D_{1}\right]^{1 / 2} \Delta t}$
So, we have that
$\Delta D_{1}=\frac{2 \tau \Delta t}{M_{M}}\left[\frac{D_{1}}{G M_{E}}\right]^{1 / 2}$
That for $\Delta t=3.1 \times 10^{7} \mathrm{~s}=1$ year, gives $\Delta D_{1}=0.034 \mathrm{~m}$.
This is the yearly increase in the Earth-Moon distance.

| 3 h | We now use that <br> $\tau=-\frac{I_{E} \Delta \omega_{E 1}}{\Delta t}$ <br>  <br> from where we get <br> $\Delta \omega_{E 1}=-\frac{\tau \Delta t}{I_{E}}$ <br> that for $\Delta t=3.1 \times 10^{7} \mathrm{~s}=1$ year gives <br> $\Delta \omega_{E 1}=-1.6 \times 10^{-14} \mathrm{~s}^{-1}$. <br> If $P_{E}$ is the period of time considered, we have that: <br> $\frac{\Delta P_{E}}{P_{E}}=-\frac{\Delta \omega_{E 1}}{\omega_{E}}$ <br> since $P_{E}=1$ day $=8.64 \times 10^{4} \mathrm{~s}$, we get <br> $\Delta P_{E}=1.9 \times 10^{-5} \mathrm{~s}$. |  |
| :--- | :--- | :--- |
| This is the amount of time that the day lengthens in a year. |  |  |

## 4. Where is the energy going?

| 4 a | The present total (rotational plus gravitational) energy of the system is: | 0.4 |
| :--- | :--- | :--- |
|  | $E=\frac{1}{2} I_{E} \omega_{E 1}^{2}+\frac{1}{2} I_{M} \omega_{M 1}^{2}-\frac{G M_{E} M_{M}}{D_{1}}$. |  |
|  | Using that |  |
|  | $\omega_{M 1}^{2} D_{1}^{3}=G M_{E}$, we get |  |

$$
E=\frac{1}{2} I_{E} \omega_{E 1}^{2}-\frac{1}{2} \frac{G M_{E} M_{M}}{D_{1}}
$$

| 4 b | $\Delta E=I_{E} \omega_{E 1} \Delta \omega_{E 1}+\frac{1}{2} \frac{G M_{E} M_{M}}{D_{1}^{2}} \Delta D_{1}$, that gives | 0.4 |
| :--- | :--- | :--- |
| $\Delta E=-9.0 \times 10^{19} \mathrm{~J}$ |  |  |


| 4 c | $M_{\text {water }}=4 \pi r_{o}^{2} \times h \times \rho_{\text {water }} \mathrm{kg}=2.6 \times 10^{17} \mathrm{~kg}$. | 0.2 |
| :--- | :--- | :--- |


| 4 d | $\Delta E_{\text {water }}=-g M_{\text {water }} \times 0.5 \mathrm{~m} \times 2$ day $^{-1} \times 365$ days $\times 0.1=-9.3 \times 10^{19}$ J. Then, the <br> two energy estimates are comparable. | 0.3 |
| :--- | :--- | :--- |

